Applications of (a, b)-continued fraction transformations

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Dedicated to the memory of Dan Rudolph

Abstract. We describe a general method of arithmetic coding of geodesics on the modular surface based on the study of one-dimensional Gauss-like maps associated to a two-parameter family of continued fractions introduced in [Katok and Ugarcovici. Structure of attractors for (a, b)-continued fraction transformations. J. Modern Dynamics 4 (2010), 637–691]. The finite rectangular structure of the attractors of the natural extension maps and the corresponding 'reduction theory' play an essential role. In special cases, when an (a, b)-expansion admits a so-called 'dual', the coding sequences are obtained by juxtaposition of the boundary expansions of the fixed points, and the set of coding sequences is a countable sofic shift. We also prove that the natural extension maps are Bernoulli shifts and compute the density of the absolutely continuous invariant measure and the measure-theoretic entropy of the one-dimensional map.

1. Introduction and background

In [16], the authors studied a new two-parameter family of continued fraction transformations. These transformations can be defined using the standard generators T(x) = x + 1, S(x) = -1/x of the modular group $SL(2, \mathbb{Z})$ and considering $f_{a,b} : \mathbb{R} \to \mathbb{R}$ given by

$$f_{a,b}(x) = \begin{cases} x+1 & \text{if } x < a, \\ -\frac{1}{x} & \text{if } a \le x < b, \\ x-1 & \text{if } x \ge b. \end{cases}$$
 (1.1)

Under the assumption that the parameters (a, b) belong to the set

$$\mathcal{P} = \{(a, b) \mid a < 0 < b, b - a > 1, -ab < 1\},\$$

one can introduce corresponding continued fraction algorithms by using the first return map of $f_{a,b}$ to the interval [a, b). Equivalently, these so-called (a, b)-continued fractions can be defined using a generalized integral part function:

$$\lfloor x \rceil_{a,b} = \begin{cases} \lfloor x - a \rfloor & \text{if } x < a, \\ 0 & \text{if } a \le x < b, \\ \lceil x - b \rceil & \text{if } x \ge b, \end{cases}$$
 (1.2)

where $\lfloor x \rfloor$ denotes the integer part of x and $\lceil x \rceil = \lfloor x \rfloor + 1$.

A starting point of the theory is the following result [16, Theorem 2.1]: if $(a, b) \in \mathcal{P}$, then any irrational number x can be expressed uniquely as an infinite continued fraction of the form

$$x = n_0 - \frac{1}{n_1 - \frac{1}{n_2 - \frac{1}{\cdots}}} = \lfloor n_0, n_1, \dots \rceil_{a,b} \quad (n_k \neq 0 \text{ for } k \geq 1),$$

where $n_0 = \lfloor x \rceil_{a,b}$, $x_1 = -1/(x - n_0)$ and $n_{k+1} = \lfloor x_{k+1} \rceil_{a,b}$, $x_{k+1} = -1/(x_k - n_k)$, i.e. the sequence of partial fractions $r_k = \lfloor n_0, n_1, \dots, n_k \rceil_{a,b}$ converges to x.

It is possible to construct (a, b)-continued fraction expansions for rational numbers, too. However, such expansions will terminate after finitely many steps if $b \neq 0$. If b = 0, the expansions of rational numbers will end with a tail of 2s, since $0 = [1, 2, 2, ...]_{a,0}$.

The above family of continued fraction transformations contains three classical examples: the case a = -1, b = 0 described in [12, 22] gives the 'minus' (backward) continued fractions, the case a = -1/2, b = 1/2 gives the 'closest-integer' continued fractions considered first by Hurwitz in [9], and the case a = -1, b = 1 was presented in [14, 19] in connection with a method of coding symbolically the geodesic flow on the modular surface following Artin's pioneering work [6] and corresponds to the regular 'plus' continued fractions with alternating signs of the digits.

The main object of study in [16] is a two-dimensional realization of the *natural* extension map of $f_{a,b}$, $F_{a,b}: \mathbb{R}^2 \setminus \Delta \to \mathbb{R}^2 \setminus \Delta$, $\Delta = \{(x, y) \in \mathbb{R}^2 \mid x = y\}$, defined by

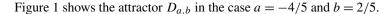
$$F_{a,b}(x, y) = \begin{cases} (x+1, y+1) & \text{if } y < a, \\ \left(-\frac{1}{x}, -\frac{1}{y}\right) & \text{if } a \le y < b, \\ (x-1, y-1) & \text{if } y \ge b. \end{cases}$$
 (1.3)

Here is the main result of that paper.

THEOREM 1.1. [16] There exists an explicit uncountable set \mathcal{E} of one-dimensional Lebesgue measure zero that lies on the diagonal boundary b = a + 1 of \mathcal{P} such that:

- (1) for all $(a, b) \in \mathcal{P} \setminus \mathcal{E}$ the map $F_{a,b}$ has an attractor $D_{a,b} = \bigcap_{n=0}^{\infty} F_{a,b}^n(\mathbb{R}^2 \setminus \Delta)$ on which $F_{a,b}$ is essentially bijective;
- (2) the set $D_{a,b}$ consists of two (or one, in degenerate cases) connected components each having finite rectangular structure, i.e. bounded by non-decreasing step-functions with a finite number of steps;

(3) almost every point (x, y) of the plane $(x \neq y)$ is mapped to $D_{a,b}$ after finitely many iterations of $F_{a,b}$.



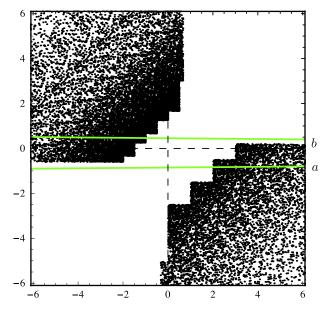


FIGURE 1. Attracting domain $D_{a,b}$ for $a = -\frac{4}{5}$, $b = \frac{2}{5}$.

An essential role in the argument is played by the forward orbits associated to a and b: to a, the *upper orbit* $\mathcal{O}_u(a)$ (i.e. the orbit of Sa) and the *lower orbit* $\mathcal{O}_\ell(a)$ (i.e. the orbit of Ta); and to b, the *upper orbit* $\mathcal{O}_u(b)$ (i.e. the orbit of $T^{-1}b$) and the *lower orbit* $\mathcal{O}_\ell(b)$ (i.e. the orbit of Sb). It was proved in [16] that if $(a, b) \in \mathcal{P} \setminus \mathcal{E}$, then $f_{a,b}$ satisfies the *finiteness condition*, i.e. for both a and b, their upper and lower orbits are either eventually periodic, or they satisfy the *cycle property*, i.e. they meet forming a cycle; more precisely, there exist $k_1, m_1, k_2, m_2 \geq 0$ such that

$$f_{a,b}^{m_1}(Sa) = f_{a,b}^{k_1}(Ta) = c_a \quad \text{(respectively, } f_{a,b}^{m_2}(T^{-1}b) = f_{a,b}^{k_2}(Sb) = c_b),$$

where c_a and c_b are the *ends of the cycles*. If the products of transformations over the upper and lower sides of the cycle are equal, the cycle property is *strong*; otherwise it is *weak*. In both cases the set $\mathcal{L}_{a,b}$ of the corresponding values is finite; ends of the cycles belong to the set $\mathcal{L}_{a,b}$ if and only if they are equal to 0, i.e. if the cycle is weak. The structure of the attractor $D_{a,b}$ is explicitly 'computed' from the finite set $\mathcal{L}_{a,b}$.

The paper is organized as follows. In $\S 2$ we give some background information about geodesic flows and their representations as special flows over symbolic dynamical systems, and define the coding map. In $\S 3$ we describe the reduction procedure for coding geodesics via (a,b)-continued fractions based on the study of the attractor of the associated natural extension map, define the corresponding cross-section set, and introduce the notion of *reduced geodesic*. In $\S 4$ we prove that the first return map to the cross-section corresponds

to a shift of the coding sequence (Theorem 4.1) and, as a consequence, show that (a, b)-continued fractions satisfy the *tail property*, i.e. two $SL(2, \mathbb{Z})$ -equivalent real numbers have the same tails in their (a, b)-continued fraction expansions. In §5 we introduce a notion of a *dual code* and show that if an (a, b)-expansion has a dual (a', b')-expansion, then the coding sequence of a reduced geodesic is obtained by juxtaposition of the (a, b)-expansion of its attracting endpoint w and the (a', b')-expansion of 1/u, where u is its repelling endpoint. We also prove that if the (a, b)-expansion admits a dual, then the set of admissible coding sequences is a sofic shift (Theorem 5.8). In §6 we derive formulas for the density of the absolutely continuous invariant measure and the measure-theoretic entropy of the one-dimensional Gauss-type maps and their natural extensions. We also prove that the natural extension maps are Bernoulli shifts. Finally, in §7 we apply results of [16] to obtain explicit formulas for invariant measure for the one-dimensional maps for some regions of the parameter set \mathcal{P} .

2. Geodesic flow on the modular surface and its representation as a special flow over a symbolic dynamical system

Let $\mathcal{H}=\{z=x+iy:y>0\}$ be the upper half-plane endowed with the hyperbolic metric, $\mathcal{F}=\{z\in\mathcal{H}:|z|\geq 1,\ |\mathrm{Re}\ z|\leq 1/2\}$ be the standard fundamental region of the modular group $\mathrm{PSL}(2,\mathbb{Z})=\mathrm{SL}(2,\mathbb{Z})/\{\pm I\}$, and $M=\mathrm{PSL}(2,\mathbb{Z})\backslash\mathcal{H}$ be the modular surface. Let $S\mathcal{H}$ denote the unit tangent bundle of \mathcal{H} . We will use the coordinates $v=(z,\zeta)$ on $S\mathcal{H}$, where $z\in\mathcal{H},\ \zeta\in\mathbb{C},\ |\zeta|=\mathrm{Im}(z)$. The quotient space $\mathrm{PSL}(2,\mathbb{Z})\backslash S\mathcal{H}$ can be identified with the unit tangent bundle of M, SM, although the structure of the fibered bundle has singularities at the elliptic fixed points (see [11, §3.6] for details). Recall that geodesics in this model are half-circles or vertical half-rays. The geodesic flow $\{\tilde{\varphi}^t\}$ on \mathcal{H} is defined as an \mathbb{R} -action on the unit tangent bundle $S\mathcal{H}$ which moves a tangent vector along the geodesic defined by this vector with unit speed. The geodesic flow $\{\tilde{\varphi}^t\}$ on \mathcal{H} descends to the *geodesic flow* $\{\varphi^t\}$ on the factor M via the canonical projection

$$\pi: S\mathcal{H} \to SM$$
 (2.1)

of the unit tangent bundles. Geodesics on M are orbits of the geodesic flow $\{\varphi^t\}$.

A *cross-section C* for the geodesic flow is a subset of the unit tangent bundle *SM* visited by (almost) every geodesic infinitely often both in the future and in the past. In other words, every $v \in C$ defines an oriented geodesic $\gamma(v)$ on M which will return to C infinitely often. The 'ceiling' function $g: C \to \mathbb{R}$ giving the *time of the first return* to C is defined as follows: if $v \in C$ and t is the time of the first return of $\gamma(v)$ to C, then g(v) = t. The map $R: C \to C$ defined by $R(v) = \varphi^{g(v)}(v)$ is called the *first return map*. Thus $\{\varphi^t\}$ can be represented as a *special flow* on the space

$$C^g = \{(v, s) : v \in C, 0 \le s \le g(v)\},\$$

given by the formula $\varphi^t(v, s) = (v, s + t)$ with the identification (v, g(v)) = (R(v), 0) (see Figure 2).

Let \mathcal{N} be a finite or countable alphabet, $\mathcal{N}^{\mathbb{Z}} = \{x = \{n_i\}_{i \in \mathbb{Z}} \mid n_i \in \mathcal{N}\}$ be the space of all bi-infinite sequences endowed with the Tikhonov (product) topology,

$$\sigma: \mathcal{N}^{\mathbb{Z}} \to \mathcal{N}^{\mathbb{Z}}$$
 defined by $(\sigma x)_i = n_{i+1}$

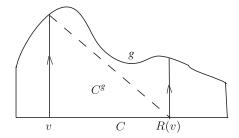


FIGURE 2. Geodesic flow is a special flow.

be the left shift map, and $\Lambda \subset \mathcal{N}^{\mathbb{Z}}$ be a closed σ -invariant subset. Then (Λ, σ) is called a *symbolic dynamical system*. There are some important classes of such dynamical systems. The space $(\mathcal{N}^{\mathbb{Z}}, \sigma)$ is called the *full shift* (or the *topological Bernoulli shift*). If the space Λ is given by a set of simple transition rules which can be described with the help of a matrix consisting of zeros and ones, we say that (Λ, σ) is a *one-step topological Markov chain* or simply a *topological Markov chain* (also called a *subshift of finite type*). A factor of a topological Markov chain is called a *sofic shift*. (See [10, §1.9] for the definitions.)

In order to represent the geodesic flow as a special flow over a symbolic dynamical system, one needs to choose an appropriate cross-section C and code it, i.e. to find an appropriate symbolic dynamical system (Λ, σ) and a continuous surjective map Cod: $\Lambda \to C$ (in some cases the actual domain of Cod is Λ except a finite or countable set of excluded sequences) defined such that the diagram

$$\begin{array}{ccc}
\Lambda & \xrightarrow{\sigma} & \Lambda \\
\text{Cod} & & \downarrow & \text{Cod} \\
& & & \downarrow & \text{Cod}
\end{array}$$

is commutative. We can then talk about *coding sequences* for geodesics defined up to a shift which corresponds to a return of the geodesic to the cross-section C. Notice that usually the coding map is not injective but only finite-to-one (see [2, §§3.2 and 5]).

There are two essentially different methods of coding geodesics on surfaces of constant negative curvature. The geometric code, with respect to a given fundamental region, is obtained by a construction universal for all Fuchsian groups. The second method is specific to the modular group and is of arithmetic nature: it uses continued fraction expansions of the end points of the geodesic at infinity and a so-called reduction theory (see [14, 15] for the three classical cases). Here we will describe a general method of arithmetic coding via (a, b)-continued fractions that is based on study of the attractor of the associated natural extension map. This approach, coupled with ideas of Bowen and Series [7], may be useful for coding of geodesics on quotients by general Fuchsian groups.

3. The reduction procedure

In what follows we will denote the end points of geodesics on \mathcal{H} by u and w, and whenever we refer to such geodesics, we use (u, w) as their coordinates on \mathbb{R}^2 $(u \neq w)$.

The reduction procedure for symbolically coding the geodesic flow on the modular surface via continued fraction expansions was presented for the three classical cases in [14]; for a survey on symbolic dynamics of the geodesic flow see also [15]. Here we describe the reduction procedure for (a, b)-continued fractions and explain how it can be used for coding purposes.

Let γ be an arbitrary geodesic on \mathcal{H} from u to w (irrational end points), and $w = \lfloor n_0, n_1, \ldots \rceil_{a,b}$. We construct the sequence of real pairs $\{(u_k, w_k)\}$ $(k \ge 0)$ defined by

$$u_0 = u, w_0 = w \text{ and } w_{k+1} = ST^{-n_k}w_k, \quad u_{k+1} = ST^{-n_k}u_k.$$
 (3.1)

Each geodesic γ_k from u_k to w_k is PSL(2, \mathbb{Z})-equivalent to γ by construction. It is convenient to describe this procedure using the *reduction map* that combines the appropriate iterate of the map $F_{a,b}$:

$$R_{a,b}: \mathbb{R}^2 \backslash \Delta \to \mathbb{R}^2 \backslash \Delta$$

given by the formula $R_{a,b}(u, w) = (ST^{-n}u, ST^{-n}w)$, where n is the first digit in the (a, b)-expansion of w. Notice that $(u_k, w_k) = R_{a,b}^k(u, w)$.

Definition 3.1. A geodesic in \mathcal{H} from u to w is called (a, b)-reduced if $(u, w) \in \Lambda_{a,b}$, where

$$\Lambda_{a,b} = F_{a,b}(D_{a,b} \cap \{a \le w \le b\}) = S(D_{a,b} \cap \{a \le w \le b\}).$$

According to Theorem 1.1, for (almost) every geodesic γ from u to w in \mathcal{H} , the above algorithm produces in finitely many steps an (a,b)-reduced geodesic PSL $(2,\mathbb{Z})$ -equivalent to γ , and an application of this algorithm to an (a,b)-reduced geodesic produces another (a,b)-reduced geodesic. In other words, there exists a positive integer ℓ such that $R_{a,b}^{\ell}(u,w) \in \Lambda_{a,b}$ and $R_{a,b}: \Lambda_{a,b} \to \Lambda_{a,b}$ is bijective (with the exception of some segments of the boundary of $\Lambda_{a,b}$ and their images).

Let γ be a reduced geodesic with the repelling point $u \neq 0$ and the attracting point

$$w = [n_0, n_1, \dots]_{a,b}.$$
 (3.2)

Then, by successive applications of the map $R_{a,b}$, we obtain a sequence of real pairs $\{(u_k, w_k)\}$ $(k \ge 0)$ (see (3.1) above) such that each geodesic γ_k from u_k to w_k is (a, b)-reduced. Using the bijectivity of the map $R_{a,b}$, we extend the sequence (3.2) to the past to obtain a bi-infinite sequence of integers

$$\lfloor \gamma \rceil = \lfloor \dots, n_{-2}, n_{-1}, n_0, n_1, n_2, \dots \rceil,$$
 (3.3)

called the *coding sequence* of γ , as follows. There exist an integer $n_{-1} \neq 0$ and a real pair $(u_{-1}, w_{-1}) \in \Lambda_{a,b}$ such that $ST^{-n_{-1}}w_{-1} = w = w_0$ and $ST^{-n_{-1}}u_{-1} = u = u_0$. Notice that $\lfloor w_{-1} \rceil_{a,b} = n_{-1}$. By uniqueness of the (a,b)-expansion, we conclude that $w_{-1} = \lfloor n_{-1}, n_0, n_1, \ldots \rceil_{a,b}$. Continuing inductively, we define the sequence of integers n_{-k} and the real pairs $(u_{-k}, w_{-k}) \in \Lambda_{a,b}$ $(k \geq 2)$, where

$$w_{-k} = [n_{-k}, n_{-k+1}, n_{-k+2}, \ldots]_{a,b},$$

by $ST^{-n_{-k}}w_{-k}=w_{-(k-1)}$ and $ST^{-n_{-k}}u_{-k}=u_{-(k-1)}$. We also associate to γ a bi-infinite sequence of (a,b)-reduced geodesics

$$(\ldots, \gamma_{-2}, \gamma_{-1}, \gamma_0, \gamma_1, \gamma_2, \ldots),$$
 (3.4)

where γ_k is the geodesic from u_k to w_k .

Remark 3.2. Notice that all 'intermediate' geodesics $T^{-s}\gamma_k$ $(1 \le s \le n_k)$ obtained from γ_k using the map $F_{a,b}$ are not (a,b)-reduced.

PROPOSITION 3.3. A formal minus continued fraction consisting of the digits of the 'past' of (3.3),

$$n_{-1} - \frac{1}{n_{-2} - \frac{1}{n_{-3} - \frac{1}{\cdots}}} = (n_{-1}, n_{-2}, n_{-3}, \dots)$$

converges to 1/u.

Proof. By [13, Lemma 1.1], it will be sufficient to check that $|n_{-k}| = 1$ implies $n_{-k} \cdot n_{-(k+1)} < 0$, i.e. the digit 1 must be followed by a negative integer and the digit -1 must be followed by a positive integer. We use the following properties of the set $\Lambda_{a,b}$ that can be derived from our knowledge of the shape of the set $D_{a,b}$ determined in [16, Lemmas 5.6, 5.10, 5.11]. The upper part of $\Lambda_{a,b}$ is contained in the region

$$[-1,0] \times \left[-\frac{1}{a}, +\infty \right] \cup [0,1] \times \left[-\frac{1}{b-1}, +\infty \right] \quad \text{if } b < 1,$$

$$[-1,0] \times \left[-\frac{1}{a}, +\infty \right] \qquad \qquad \text{if } b \ge 1.$$

$$(3.5)$$

The lower part of $\Lambda_{a,b}$ is contained in the region

$$[-1,0] \times \left[-\infty, -\frac{1}{a+1}\right] \cup [0,1] \times \left[-\infty, -\frac{1}{b}\right] \quad \text{if } a > -1,$$

$$[0,1] \times \left[-\infty, -\frac{1}{b}\right] \qquad \qquad \text{if } a \leq -1.$$

$$(3.6)$$

Recall that $(u_{-(k+1)}, w_{-(k+1)}) = (T^{n_{-(k+1)}}Su_{-k}, T^{n_{-(k+1)}}Sw_{-k})$ for an appropriate integer $n_{-(k+1)} \neq 0$. Suppose that $n_{-k} = 1$. Then $w_{-k} > 0$. If $u_{-k} < 0$, then $Su_{-k} > 0$ and $Sw_{-k} < 0$, and it takes a negative power of T to bring it back to (the lower component of) $\Lambda_{a,b}$, i.e. $n_{-(k+1)} < 0$. The case $u_{-k} > 0$, according to (3.5), can only occur if $b \leq 1$. In this case, $-1/(b-1) \leq w_{-k} < b+1$, which is equivalent to b > 1, a contradiction. Therefore $n_{-k} = 1$ implies $n_{-(k+1)} < 0$. A similar argument shows that $n_{-k} = -1$ implies $n_{-(k+1)} > 0$. We conclude that the formal minus continued fraction converges. In order to prove that the limit is equal to 1/u we use the recursive definition of the digits n_{-1}, n_{-2}, \ldots , to write

$$\frac{1}{u} = n_{-1} - u_{-1} = n_{-1} - \frac{1}{n_{-2} - u_{-2}} = \dots = (n_{-1}, n_{-2}, \dots, n_{-k} - u_{-k}) = \dots,$$

and the conclusion follows since the formal minus continued fraction converges. \Box

Let

$$C = \{z \in \mathcal{H} \mid |z| = 1, -1 \le \text{Re } z \le 1\}$$

be the upper-half of the unit circle, and

$$C^{-} = \{z \in \mathcal{H} \mid |z+1| = 1, -\frac{1}{2} \le \operatorname{Re} z \le 0\}$$

and

$$C^+ = \{ z \in \mathcal{H} \mid |z - 1| = 1, \ 0 \le \operatorname{Re} z \le \frac{1}{2} \}$$

be the images of the two vertical boundary components of the fundamental region \mathcal{F} under S (see Figure 3).

PROPOSITION 3.4. Every (a, b)-reduced geodesic intersects either C or both curves C^- and C^+ .

Proof. If a, b are such that $-1 \le a \le 0$ and $0 \le b \le 1$, then by properties (3.5) and (3.6) of the set $\Lambda_{a,b}$, if $(u, w) \in \Lambda_{a,b}$, then $-1 \le u \le 1$ and $w \ge -1/a$ or $w \le -1/b$, and hence all (a, b)-reduced geodesics intersect C. For the case b > 1 we have: if -1 < u < 0, then either w > -1/a > b > 1 or w < -1/(a+1) < -1, i.e. the geodesic intersects C; if 0 < u < 1, then (3.5) implies that w < -1/b < a < 0, thus the corresponding geodesic intersects C if w < -1, and it intersects first C^+ and then C^- if -1 < w < 0. Similarly, for the case a < -1 we have: if 0 < u < 1, then either w < -1/b < a < -1 or w > -1/(b-1) > 1, i.e. the geodesic intersects C; if -1 < u < 0, then (3.6) implies that w > -1/a > b > 0, therefore the corresponding geodesic intersects C if w > 1, and it intersects first C^- and then C^+ if 0 < w < 1. □

Based on Proposition 3.4 we introduce the notion of the *cross-section point*. It is either the intersection of a reduced geodesic γ with C, or, if γ does not intersect C, its first intersection with $C^- \cup C^+$.

Now we can define a map

$$\varphi: \Lambda_{a,b} \to S\mathcal{H}, \quad \varphi(u, w) = (z, \zeta),$$

where $z \in \mathcal{H}$ is the cross-section point on the geodesic γ from u to w, and ζ is the unit vector tangent to γ at z. The map is clearly injective. Composed with the canonical projection π introduced in (2.1), we obtain a map

$$\pi \circ \varphi : \Lambda_{a,b} \to SM$$
.

Let $C_{a,b} = \pi \circ \varphi(\Lambda_{a,b}) \subset SM$. This set can be described as follows: $C_{a,b} = P \cup Q_1 \cup Q_2$, where P consists of the unit vectors based on the circular boundary of the fundamental region \mathcal{F} pointing inward such that the corresponding geodesic γ on the upper half-plane \mathcal{H} is (a,b)-reduced, Q_1 consists of the unit vectors based on the right vertical boundary of \mathcal{F} pointing inward such that either $S\gamma$ or $TS\gamma$ is (a,b)-reduced (notice that they cannot both be reduced), and Q_2 consists of the unit vectors based on the left vertical boundary of \mathcal{F} pointing inward such that either $S\gamma$ or $T^{-1}S\gamma$ is (a,b)-reduced (see Figure 3). Then almost every orbit of $\{\varphi^t\}$ returns to $C_{a,b}$, i.e. $C_{a,b}$ is a *cross-section* for $\{\varphi^t\}$, and $\Lambda_{a,b}$ is a parametrization of $C_{a,b}$. The map $\pi \circ \varphi$ is injective, as follows from Remark 3.2: only one of the geodesics γ , $S\gamma$, $T^{-1}S\gamma$, and $TS\gamma$ can be reduced.

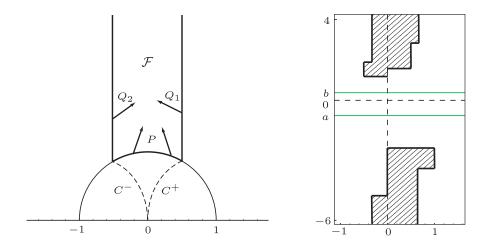


FIGURE 3. The cross-section (left) and its $\Lambda_{a,b}$ parametrization (right).

4. Symbolic coding of the geodesic flow via (a, b)-continued fractions If γ is a geodesic on \mathcal{H} , we denote by $\bar{\gamma}$ the canonical projection of γ on M. For a given geodesic on M that can be reduced in finitely many steps, we can always choose its lift γ to \mathcal{H} to be (a, b)-reduced.

The following theorem provides the basis for coding geodesics on the modular surface using (a, b)-coding sequences.

Theorem 4.1. Let γ be an (a,b)-reduced geodesic on $\mathcal H$ and $\bar{\gamma}$ its projection to M. Then:

- (1) each geodesic segment of $\bar{\gamma}$ between successive returns to the cross-section $C_{a,b}$ produces an (a, b)-reduced geodesic on \mathcal{H} , and each reduced geodesic $SL(2, \mathbb{Z})$ -equivalent to γ is obtained in this way;
- (2) the first return of $\bar{\gamma}$ to the cross-section $C_{a,b}$ corresponds to a left shift of the coding sequence of γ .

Proof. (1) By lifting a geodesic segment on M starting on $C_{a,b}$ to \mathcal{H} , we obtain a segment of a geodesic γ on \mathcal{H} that is reduced by the definition of the cross-section $C_{a,b}$. A coding sequence of $\gamma = \gamma_0$ that connects u_0 to $w_0 = \lfloor n_0, n_1, \ldots \rceil_{a,b}$,

$$[\gamma_0] = [\ldots, n_{-2}, n_{-1}, n_0, n_1, n_2, \ldots],$$

is obtained by extending the sequence of digits of w_0 to the past as explained in the previous section.

Let us assume that $w_0 > 0$, i.e. $n_0 \ge 1$. The case $w_0 < 0$ can be treated similarly. The geodesic $ST^{-n_0}\gamma_0 = \gamma_1$ is reduced by Theorem 1.1. Let z_0 and z_1 be the cross-section points on γ_0 and γ_1 , respectively. Then $z_1' = T^{n_0}Sz_1 \in \gamma_0$; it is the intersection point of γ_0 with the circle $|z - n_0| = 1$. We will show that the geodesic segment of γ_0 , $[z_0, z_1']$, projected to M is the segment between two successive returns to the cross-section $C_{a,b}$. Since $ST^{-n_0}(z_1') = z_1$ is the cross-section point on γ_1 , the geodesic segment

 $[z_0, z_1']$ projected to M is between two returns to $C_{a,b}$. Recall that a geodesic in \mathcal{F} consists of countably many oriented geodesic segments between consecutive crossings of the boundary of \mathcal{F} that are obtained by the canonical projection of γ_0 to \mathcal{F} .

If z_0 is the intersection of γ_0 with C, there are two possibilities: first, when γ_0 intersects \mathcal{F} or γ_0 does not intersect \mathcal{F} and $ST^{-1}\gamma_0$ exits \mathcal{F} through its circular boundary; and second, when γ_0 does not intersect \mathcal{F} and $ST^{-1}\gamma_0$ exits \mathcal{F} through its left vertical boundary. In the first case the segments in \mathcal{F} are represented by the intersection with \mathcal{F} of the following geodesics in $\mathcal{H}: T^{-1}\gamma_0, T^{-2}\gamma_0, \ldots, T^{-n_0+1}\gamma_0$, either $ST^{-n_0+1}\gamma_0$ or $T^{-n_0}\gamma_0$, and either γ_0 or $ST^{-1}\gamma_0$.

Suppose that for some intermediate point $z \in \gamma_0$, $z \in [z_0, z_1']$ the unit vector tangent to γ_0 at z, (z, ζ) , is projected to $C_{a,b}$. By tracing the geodesic γ_0 inside \mathcal{F} , we see that (z, ζ) must be projected to $(\bar{z}, \bar{\zeta})$ with \bar{z} on the boundary of \mathcal{F} and $\bar{\zeta}$ directed inward. Then the geodesic through $(\bar{z}, \bar{\zeta})$:

- (a) enters \mathcal{F} through its vertical boundary and exits it also through the vertical boundary;
- (b) enters \mathcal{F} through its vertical boundary and exits through its circular boundary; or
- (c) enters \mathcal{F} through its circular boundary and exits through its vertical boundary. The following assertions are implied by the analysis of the attractor $D_{a,b}$. In case (a), $T^{-1}ST^{-s}\gamma_0$ is not reduced for $1 \le s < n_0$ since $s < n_0$, $T^{-s}w_0 > b$, hence $ST^{-s}w_0 > -1/b$, i.e. $(ST^{-s}u_0, ST^{-s}w_0) \notin D_{a,b}$, therefore

$$(T^{-1}ST^{-s}u_0, T^{-1}ST^{-s}w_0) \notin \Lambda_{a,b}.$$

In case (b), if the segment $T^{-n_0}\gamma_0$ exits through the circular boundary of \mathcal{F} , $ST^{-n_0}\gamma_0 = \gamma_1$ is reduced, and we reached the point z_1 on the cross-section. If the segment $T^{-n_0+1}\gamma_0$ intersects the circular boundary of \mathcal{F} , $ST^{-n_0+1}\gamma_0$ is not reduced. In case (c), ST^{-n_0+1} is not reduced.

In the second case the first digit of w_0 is $n_0 = 2$. This is because $n_0 = 1$ would imply that b + 1 < w < -1/(b-1) which is impossible. Thus $ST^{-2}\gamma_0 = \gamma_1$ is reduced. In this case the geodesic in \mathcal{F} consists of the intersection with \mathcal{F} of a single geodesic $ST^{-1}\gamma_0$ that enters \mathcal{F} through its right vertical and leaves it through its left vertical boundary, since $(TS)T(ST^{-1}\gamma_0) = ST^{-2}\gamma_0 = \gamma_1$ is reduced. In all cases the geodesic segment $[z_0, z_1']$ projected to M is between two consecutive returns to $C_{a,b}$.

If $z_0 \notin C$, by Proposition 3.4, since $w_0 > 0$, $z_0 \in C^-$. Notice that this implies that a < -1 and $n_0 = 1$, and $\gamma_1 = ST^{-1}\gamma_0$ is reduced. In this case the geodesic in \mathcal{F} also consists of the intersection with \mathcal{F} of a single geodesic $S\gamma_0$ that enters \mathcal{F} through its right vertical and leaves it through its left vertical boundary, since $(TS)T(S\gamma_0) = ST^{-1}\gamma_0 = \gamma_1$ is reduced, and hence the geodesic segment $[z_0, z_1']$ projected to M is between two consecutive returns to $C_{a,b}$. Continuing this argument by induction in both the positive and negative direction, we obtain a bi-infinite sequence of points

$$(\ldots, z_{-2}, z_{-1}, z_0, z_1, z_2, \ldots),$$

where z_k is the cross-section point of the reduced geodesic γ_k in the sequence of γ_0 , that represents the sequence of all successive returns of the geodesic γ_0 in M to the cross-section $C_{a,b}$.

If $\tilde{\gamma}_0$ is a reduced geodesic in \mathcal{H} , $SL(2, \mathbb{Z})$ -equivalent to γ_0 , then both project to the same geodesic on M. Therefore, the cross-section point \tilde{z}_0 of $\tilde{\gamma}_0$ projects on $C_{a,b}$ to a cross-section point z_k of γ_k for some k. This completes the proof of (1).

(2) Since $\gamma_1 = ST^{-n_0}\gamma_0$, $w_1 = ST^{-n_0}w_0 = \lfloor n_1, n_2, \ldots \rceil_{a,b}$. The first digit of the past is evidently n_0 , and the remaining digits are the same as for γ_0 . Thus (2) follows.

The following corollary is immediate.

COROLLARY 4.2. If γ' is $SL(2, \mathbb{Z})$ -equivalent to γ , and both geodesics can be reduced in finitely many steps, then the coding sequences of γ and γ' differ by a shift.

This implies a very important property of (a, b)-continued fractions that escapes a direct proof.

COROLLARY 4.3. (The tail property) For almost every pair of real numbers that are $SL(2, \mathbb{Z})$ -equivalent, the 'tails' of their (a, b)-continued fraction expansions coincide.

Remark 4.4. The set of exceptions in Corollary 4.3 is the same as the set described in Theorem 1.1(3).

Thus we can talk about *coding sequences of geodesics on M*. To any geodesic γ that can be reduced in finitely many steps we associate the coding sequence (3.3) of a reduced geodesic $SL(2, \mathbb{Z})$ -equivalent to it. Corollary 4.2 implies that this definition does not depend on the choice of a particular representative: sequences for equivalent reduced geodesics differ by a shift.

Let $X_{a,b}$ be the closure of the set of admissible sequences and σ be the left shift map. The coding map Cod: $X_{a,b} \to C_{a,b}$ is defined by

$$Cod(\lfloor \dots, n_{-2}, n_{-1}, n_0, n_1, \dots \rceil) = (1/(n_{-1}, n_{-2}, \dots), \lfloor n_0, n_1, \dots \rceil_{a,b}).$$
 (4.1)

This map is essentially bijective.

The symbolic system $(X_{a,b}, \sigma) \subset (\mathcal{N}^{\mathbb{Z}}, \sigma)$ is defined on the infinite alphabet $\mathcal{N} \subset \mathbb{Z}\setminus\{0\}$. The product topology on $\mathcal{N}^{\mathbb{Z}}$ is induced by the distance function

$$d(x, x') = \frac{1}{m},$$

where $x = (n_i)$, $x' = (n'_i) \in \mathcal{N}^{\mathbb{Z}}$, and $m = \max\{k \mid n_i = n'_i \text{ for } |i| \le k\}$.

PROPOSITION 4.5. The map Cod is continuous.

Proof. If d(x, x') < 1/m, then the (a, b)-expansions of the attracting end points w(x) and w(x') of the corresponding geodesics given by (3.2) have the same first m digits. Hence the first m convergents of their (a, b)-expansions are the same, and, using the properties of (a, b)-continued fractions and the rate of convergence of [16, Theorem 2.1], we obtain |w(x) - w(x')| < 2/m. Similarly, the first m digits in the convergent formal minus continued fraction of 1/u(x) and 1/u(x') are the same, and hence |u(x) - u(x')| < 2|u(x)u'(x)|/m < 2/m. Therefore the geodesics are uniformly 2/m-close. But the tangent vectors v(x), $v(x') ∈ C_{a,b}$ are determined by the intersection of the corresponding geodesic with the unit circle or the curves C^+ and C^- . Hence, by making m large enough we can make v(x') as close to v(x) as we wish. □

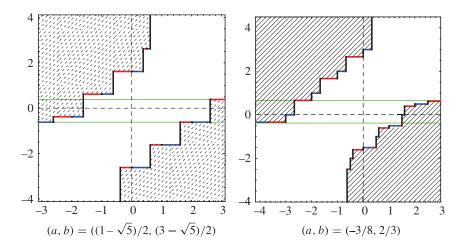


FIGURE 4. Domains of self-dual expansions.

In conclusion, the geodesic flow becomes a special flow over a symbolic dynamical system $(X_{a,b}, \sigma)$ on the infinite alphabet $\mathcal{N} \subset \mathbb{Z}\setminus\{0\}$. The ceiling function $g_{a,b}(x)$ on $X_{a,b}$ coincides with the time of the first return of the associated geodesic $\gamma(x)$ to the cross-section $C_{a,b}$. One can establish an explicit formula for $g_{a,b}(x)$ as the function of the end points of the corresponding geodesic $\gamma(x)$, u(x), w(x), following the ideas explained in [8]. If $-1 \le a \le 0$ and $0 \le b \le 1$, then $g_{a,b}(x)$ is cohomologous to $2 \log |w(x)|$; more precisely,

$$g_{a,b}(x) = 2\log|w(x)| + \log h(x) - \log h(\sigma x)$$

where

$$h(x) = \frac{|w(x) - u(x)|\sqrt{w(x)^2 - 1}}{w(x)^2\sqrt{1 - u(x)^2}}.$$

Dual codes

We have seen that a coding sequence for a reduced geodesic from u to w (see (3.3)) is comprised of the sequence of digits in the (a, b)-expansion of w and the 'past', an infinite sequence of non-zero integers, each digit of which depends on w and u. In some special cases the 'past' only depends on u, and, in fact, will coincide with the sequence of digits of 1/u by using a so-called *dual expansion* to (a, b).

Let $\psi(x, y) = (-y, -x)$ be the reflection of the plane about the line y = -x.

Definition 5.1. If $\psi(D_{a,b})$ coincides with the attractor set $D_{a',b'}$ for some $(a',b') \in \mathcal{P}$, then the (a',b')-expansion is called the *dual* expansion to (a,b). If (a',b')=(a,b), then the (a,b)-expansion is called *self-dual*.

Example 5.2. The classical situations of (-1, 0)- and (-1, 1)-expansions are self-dual. Two more sophisticated examples, $((1 - \sqrt{5})/2, (3 - \sqrt{5})/2)$ and (-3/8, 2/3), are shown in Figure 4.

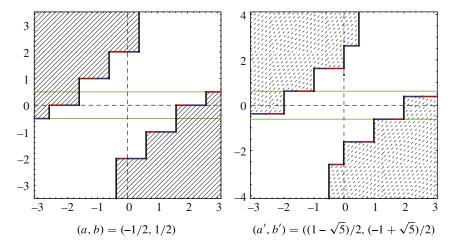


FIGURE 5. Dual expansions.

Example 5.3. The expansions $(-1/n, 1 - (1/n)), n \ge 1$, satisfy a weak cycle property and have dual expansions that are periodic. A classical example in this series is the Hurwitz case (-1/2, 1/2) whose dual is $((1 - \sqrt{5})/2, (-1 + \sqrt{5})/2)$ (see [9, 14]). Their domains are shown in Figure 5.

The following result gives equivalent characterizations for an expansion to admit a dual.

PROPOSITION 5.4. The following are equivalent:

- (i) the (a, b)-expansion has a dual;
- (ii) the boundary of the lower part of the set $D_{a,b}$ does not have y-levels with a < y < 0, and the boundary of the upper part of the set $D_{a,b}$ does not have y-levels with 0 < y < b;
- (iii) a and b do not have the strong cycle property.

Proof. If the (a, b)-expansion has a dual (a', b')-expansion, then the parameters a', b' are obtained from the boundary of $D_{a,b}$ as follows: the right vertical boundary of the upper part of $D_{a,b}$ is the ray x = 1 - b', and the left vertical boundary of the lower part of $D_{a,b}$ is the ray x = -1 - a'. Now assume that (ii) does not hold. Then at least one of the parameters a, b has the strong cycle property, and either the left boundary of the upper part of $\Lambda_{a,b}$ or the right boundary of the lower part of $\Lambda_{a,b}$ is not a straight line. Assume the former. Then the reflection of $D_{a,b}$ with respect to the line y = -x is not $D_{a',b'}$ since the map $F_{a',b'}$ is not bijective on it: the black rectangle in Figure 6 belongs to it, but its image under T^{-1} , colored in grey, does not. Thus (i) \Rightarrow (ii).

Conversely, let the vertical line x = 1 - b' be the right boundary of the upper part of $D_{a,b}$ and the vertical line x = -1 - a' be the left boundary of the lower part of $D_{a,b}$. Let $[x_a, \infty] \times \{a\}$ be the intersection of $D_{a,b}$ with the horizontal line at the level a, and $[-\infty, x_b] \times \{b\}$ be the intersection of $D_{a,b}$ with the horizontal line at the level b. Then $a' = 1/x_b$ and $b' = 1/x_a$. We also see that 1 - b' = -1/t, where $t = x_b$ or $t < x_b$ if

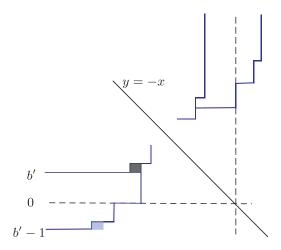


FIGURE 6. Dual expansions and $D_{a,b}$.

 $[t, x_b] \times \{0\}$ is a segment of the boundary of $D_{a,b}$. Then $-b' + 1 = -1/t \le a'$, which implies that $b' - a' \ge 1$. By [16, Lemma 5.6], $x_b \le -1$ and $x_a \ge 1$, therefore

$$-1 \le a' \le 0 \le b' \le 1,\tag{5.1}$$

and

$$\Lambda_{a,b} = D_{a,b} \cap \{(u, w) \in \mathbb{R}^2 : -b' \le u \le -a'\}. \tag{5.2}$$

We now show that $\psi(D_{a,b}) = D_{a',b'}$ is the attractor for $F_{a',b'}$, where

$$F_{a',b'} = \psi \circ F_{a,b}^{-1} \circ \psi^{-1}.$$
 (5.3)

For $(u, w) \in D_{a',b'}$ with a' < w < b', we have $\psi^{-1}(u, w) = (-w, -u)$ with -b' < u < -a', so $\psi^{-1}(u, w) \in \Lambda_{a,b}$ by (5.2), hence $F_{a,b}^{-1}(-w, -u) = (1/w, 1/u)$, and $F_{a',b'}(u, w) = (-1/u, -1/w)$. For $(u, w) \in D_{a',b'}$ with w > b', we have $\psi^{-1}(u, w) = (-w, -u)$ with u < -b', so $F_{a,b}^{-1}(-w, -u) = (-w + 1, -u + 1)$, and $F_{a',b'}(u, w) = (u - 1, w - 1)$. Similarly, for $(u, w) \in D_{a',b'}$ with w < a', we have $\psi^{-1}(u, w) = (-w, -u)$ with u > -a', so $F_{a,b}^{-1}(-w, -u) = (-w - 1, -u - 1)$, and $F_{a',b'}(u, w) = (u + 1, w + 1)$. This proves that (ii) \Rightarrow (i).

Notice that (ii) and (iii) are equivalent by [16, Theorems 4.2 and 4.5].

Remark 5.5. Notice that if an (a, b)-expansion has a dual, then $-1 \le a \le 0 \le b \le 1$. This follows from (5.1) and the fact that the relation of duality is symmetric.

THEOREM 5.6. If an (a, b)-expansion admits a dual expansion (a', b'), and γ_0 is an (a, b)-reduced geodesic, then its coding sequence,

$$|\gamma_0| = |\dots, n_{-2}, n_{-1}, n_0, n_1, n_2, \dots|,$$
 (5.4)

is obtained by juxtaposing the (a, b)-expansion of $w_0 = \lfloor n_0, n_1, n_2, \ldots \rceil_{a,b}$ and the (a', b')-expansion of $1/u_0 = \lfloor n_{-1}, n_{-2}, \ldots \rceil_{a',b'}$. This property is preserved under the left shift of the sequence.

Proof. We will show that the digits in the (a', b')-expansion of $1/u_0$ coincide with the digits of the 'past' of (5.4). By (5.3), the diagram

$$\begin{array}{c|c} \Lambda_{a,b} & \xrightarrow{S\psi} & \Lambda_{a',b'} \\ R_{a,b}^{-1} & & & \downarrow R_{a',b'} \\ \Lambda_{a,b} & \xrightarrow{S\psi} & \Lambda_{a',b'} \end{array}$$

is commutative. The pair $(u_0, w_0) \in \Lambda_{a,b}$, therefore $(Su_0, Sw_0) \in S\Lambda_{a,b} \subset D_{a,b}$, and $(1/w_0, 1/u_0) \in \Lambda_{a',b'}$. The first digit of the (a', b')-expansion of $1/u_0$ is n_{-1} , so

$$R_{a',b'}(1/w_0, 1/u_0) = (ST^{-n_{-1}}(1/w_0), ST^{-n_{-1}}(1/u_0))$$

maps $\Lambda_{a',b'}$ to itself. Then

$$(u_{-1}, w_{-1}) := R_{a,b}^{-1}(u_0, w_0) = (T^{n-1}Su_0, T^{n-1}Sw_0) \in \Lambda_{a,b}$$

and

$$(ST^{-n_{-1}}u_{-1}, ST^{-n_{-1}}w_{-1}) = (u_0, w_0).$$

Also $w_{-1} = \lfloor n_{-1}, n_0, n_1, \ldots \rceil_{a,b}$ and $ST^{-n_{-1}}(1/u_0) = 1/u_{-1} = \lfloor n_{-2}, \ldots \rceil_{a',b'}$.

Continuing by induction, one proves that all digits of the 'past' of the sequence (5.4) are the digits of the (a', b')-expansion of $1/u_0$.

In order to see what happens under a left shift, we reverse the diagram to obtain

$$\begin{array}{c|c}
\Lambda_{a,b} & \xrightarrow{S\psi} & \Lambda_{a',b'} \\
R_{a,b} & & & \downarrow R_{a',b'}^{-1} \\
\Lambda_{a,b} & \xrightarrow{S\psi} & \Lambda_{a',b'}
\end{array}$$

Since the first digit of (a, b)-expansion of w_0 is n_0 ,

$$R_{a,b}(u_0, w_0) = (ST^{-n_0}u_0, ST^{-n_0}w_0)$$

maps $\Lambda_{a,b}$ to itself. Then $(u_1, w_1) := (ST^{-n_0}u_0, ST^{-n_0}w_0)$ and $w_1 = \lfloor n_1, n_2, \ldots \rceil_{a,b}$. Also

$$(1/w_1, 1/u_1) = R_{a',b'}^{-1}(1/w_0, 1/u_0) = (T_0^n S(1/w_0), T_0^n S(1/u_0)),$$

hence
$$1/u_1 = [n_0, n_{-1}, n_{-2}, \ldots]_{a',b'}$$
.

Remark 5.7. Under conditions of Theorem 5.6, if γ_0 projects to a closed geodesic on M, then its coding sequence is periodic, and $w_0 = \lfloor \overline{n_0, n_1, \dots, n_m} \rfloor_{a,b}$, $1/u_0 = \lfloor \overline{n_m, \dots, n_1, n_0} \rfloor_{a',b'}$.

THEOREM 5.8. If an (a, b)-expansion admits a dual expansion, then the symbolic space $(X_{a,b}, \sigma)$ is a sofic shift.

Proof. The 'natural' (topological) partition of the set $\Lambda_{a,b}$ related to the alphabet \mathcal{N} is $\Lambda_{a,b} = \bigcup_{n \in \mathcal{N}} \Lambda_n$, where Λ_n are labeled by the symbols of the alphabet \mathcal{N} and are defined by the condition $\Lambda_n = \{(u, w) \in \Lambda_{a,b} \mid n_0(u, w) = n_0(w) = n\}$. In order to prove that the

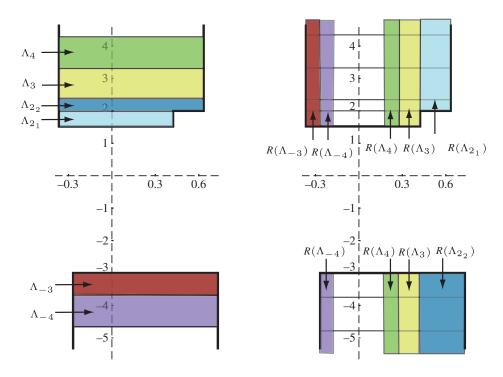


FIGURE 7. The partition of $\Lambda_{a,b}$ and its image through $R_{a,b}$.

space $(X_{a,b}, \sigma)$ is sofic one needs to find a topological Markov chain $(M_{a,b}, \tau)$ and a surjective continuous map $h: M_{a,b} \to X_{a,b}$ such that $h \circ \tau = \sigma \circ h$.

Notice that the elements Λ_n are rectangles for large n; in fact, at most two elements in the upper part and at most two elements in the lower part of $\Lambda_{a,b}$ are incomplete rectangles (see Figure 7).

Since $\Lambda_{a,b}$ has finite rectangular structure, we can subdivide these incomplete rectangles horizontally into rectangles, and extend the alphabet \mathcal{N} by adding subscripts to the corresponding elements of \mathcal{N} . For example, if Λ_2 is subdivided into two rectangles, $\Lambda_2 = \bigcup_{i=1}^2 \Lambda_{2_i}$, the 'digit' 2 will give rise to two digits, $\lambda_2 = \lambda_2 =$

We will prove that $(M_{a,b}, \tau)$ is a topological Markov chain. For this, in accordance to [2, Theorem 7.9], it is sufficient to prove that for any pair of distinct symbols $n, m \in \mathcal{N}'$, $R(M_n)$ and M_m either do not intersect or intersect 'transversally', i.e. their intersection is a rectangle with two horizontal sides belonging to the horizontal boundary of M_m

and two vertical sides belonging to the vertical boundary of $R(M_n)$. Let us recall that $-1 \le a \le 0 \le b \le 1$ (see Remark 5.5). Therefore, if $M_n = \Lambda_n$ is a complete rectangle, it is, in fact, a 1×1 square, and its image under R is an infinite vertical rectangle intersecting all M_m transversally. If M_n is obtained by subdivision of some Λ_k and belongs to the lower part of $\Lambda_{a,b}$, its horizontal boundaries are the levels of the step-function defining the lower component of $D_{a,b}$, and by Proposition 5.4, since the lower boundary of $D_{a,b}$ does not have y-levels with a < y < 0, its image is a vertical rectangle intersecting only the lower component of $D_{a,b}$ whose horizontal boundaries are the levels of the step-function defining the lower component of $D_{a,b}$. Therefore, all possible intersections with M_m are transversal. A similar argument applies to the case where M_n belongs to the upper part of $\Lambda_{a,b}$. The map $h: M_{a,b} \to X_{a,b}$ is obviously continuous, surjective, and, in addition, $h \circ \tau = \sigma \circ h$.

6. Invariant measures and ergodic properties

Based on the finite rectangular geometric structure of the domain $D_{a,b}$ and the connections with the geodesic flow on the modular surface, we study some of the measure-theoretic properties of the Gauss-type map $\hat{f}_{a,b}$: $[a,b) \rightarrow [a,b)$,

$$\hat{f}_{a,b}(x) = -\frac{1}{x} - \left[-\frac{1}{x} \right]_{a,b}, \quad \hat{f}_{a,b}(0) = 0.$$
 (6.1)

Notice that the associated natural extension map $\hat{F}_{a,b}$,

$$\hat{F}_{a,b}(x,y) = \left(\hat{f}_{a,b}(x), -\frac{1}{y - |-1/x|_{a,b}}\right),\tag{6.2}$$

is obtained from the map $F_{a,b}$ induced on the set $\Lambda_{a,b}$ by the change of coordinates

$$x = -1/w, \quad y = u \tag{6.3}$$

(or, equivalently, on the set $D_{a,b} \cap \{(u, w) \mid a \leq w < b\}$ by the change of coordinates x = w, y = -1/u). Therefore the domain $\hat{\Lambda}_{a,b}$ of $\hat{F}_{a,b}$ is easily identified knowing $\Lambda_{a,b}$ and may be considered as its 'compactification'.

Many of the measure-theoretic properties of $\hat{f}_{a,b}$ and $\hat{F}_{a,b}$ (existence of an absolutely continuous invariant measure, ergodicity) follow from the fact that the geodesic flow φ^t on the modular surface M can be represented as a special flow $(R_{a,b}, \Lambda_{a,b}, g_{a,b})$ on the space

$$\Lambda_{a,b}^{g_{a,b}} = \{(u, w, t) : (u, w) \in \Lambda_{a,b}, 0 \le t \le g_{a,b}(u, w)\}$$

(see §2). We recall that $R_{a,b} = F_{a,b}|_{\Lambda_{a,b}}$ and $g_{a,b}$ is the ceiling function (the time of the first return to the cross-section $C_{a,b}$) parametrized by $(u, w) \in \Lambda_{a,b}$.

We start with the fact that the geodesic flow $\{\varphi^t\}$ preserves the smooth (Liouville) measure $dm = du \ dw \ dt/(w-u)^2$ (see [3]), hence $R_{a,b}$ preserves the absolutely continuous measure $d\rho = du \ dw/(w-u)^2$. Using the change of coordinates (6.3), the map $\hat{F}_{a,b}$ preserves the absolutely continuous measure $dv = dx \ dy/(1+xy)^2$.

The set $\Lambda_{a,b}$ has finite measure $d\rho$ if $a \neq 0$ and $b \neq 0$, since it is uniformly bounded away from the line $\Delta = \{(u, w) : u = w\} \subset \mathbb{R}^2$ (see relations (3.5) and (3.6)). In this

situation, we can normalize the measure $d\rho$ to obtain the smooth probability measure

$$d\rho_{a,b} = \frac{d\rho}{K_{a,b}} = \frac{du \, dw}{K_{a,b}(w-u)^2} \tag{6.4}$$

where $K_{a,b} = \rho(\Lambda_{a,b})$. Similarly, if $a \neq 0$ and $b \neq 0$, the map $\hat{F}_{a,b}$ preserves the smooth probability measure

$$dv_{a,b} = \frac{dx \, dy}{K_{a,b}(1+xy)^2} \tag{6.5}$$

and $K_{a,b} = \rho(\Lambda_{a,b}) = \nu(\hat{\Lambda}_{a,b}).$

Returning to the Gauss-type map, $\hat{f}_{a,b}$, one can obtain explicitly a Lebesgue-equivalent invariant probability measure $\mu_{a,b}$ by projecting the measure $\nu_{a,b}$ onto the x-coordinate (pushforward); this is equivalent to integrating $\nu_{a,b}$ over $\hat{\Lambda}_{a,b}$ with respect to the y-coordinate as explained in [4].

We can immediately conclude that the systems $(\hat{F}_{a,b}, \nu_{a,b})$ and $(\hat{f}_{a,b}, \mu_{a,b})$ are ergodic from the fact that the geodesic flow $\{\varphi^t\}$ is ergodic with respect to dm. By using some well-known results about one-dimensional maps that are piecewise monotone and expanding, and the implications for their natural extension maps, we can establish stronger measure-theoretic properties: $(\hat{f}_{a,b}, \mu_{a,b})$ is exact and $(\hat{F}_{a,b}, \nu_{a,b})$ is a Bernoulli shift. Here we follow the presentation from [23] based on [18, 21].

THEOREM 6.1. For any $a \neq 0$ and $b \neq 0$, the system $(\hat{f}_{a,b}, \mu_{a,b})$ is exact and its natural extension $(\hat{F}_{a,b}, \nu_{a,b})$ is a Bernoulli shift.

Proof. Let us consider first the case -1 < a < 0 < b < 1. The interval (a, b) admits a countable partition $\xi = \{X_i\}_{i \in \mathbb{Z} \setminus \{0\}}$ of open intervals and the map $\hat{f}_{a,b}$ satisfies conditions (A), (F), (U) listed in [23]. Condition (A) is Adler's distortion estimate:

$$\hat{f}''_{a,b}/(\hat{f}'_{a,b})^2$$
 is bounded on $X = \bigcup_{i \in \mathbb{Z} \setminus \{0\}} X_i$, (A)

condition (F) requires the finite image property of the partition ξ ,

$$\hat{f}_{a,b}(\xi) = \{\hat{f}_{a,b}(X_i)\}_{i \in \mathbb{Z} \setminus \{0\}} \text{ is finite}, \tag{F}$$

while condition (U) is a uniformly expanding condition

$$|\hat{f}_{a,b}'| \ge \tau > 1 \text{ on } X. \tag{U}$$

Let $m \ge 0$ and $n \ge 0$ be such that $a - m - 1 \le -1/b < a - m$ and $b + n \le -1/a < b + n + 1$. Consider the open intervals

$$X_1 = \left(-\frac{1}{a-m-1}, b\right), \quad X_i = \left(-\frac{1}{a-m-i}, -\frac{1}{a-m-i+1}\right) \text{ for } i \ge 2$$

and

$$X_{-1} = \left(a, \, -\frac{1}{b+n+1}\right), \quad X_{-i} = \left(-\frac{1}{b+n+i-1}, \, -\frac{1}{b+n+i}\right) \quad \text{for } i \geq 2.$$

The map $\hat{f}_{a,b}$ satisfies conditions (A), (F), (U) with respect to the partition $\xi = \{X_i\}_{i \in \mathbb{Z} \setminus \{0\}}$. Indeed, $|\hat{f}''_{a,b}/(\hat{f}'_{a,b})^2| \leq 2$ on X, the collection of images $\hat{f}_{a,b}(\xi)$ consists of

four sets $\hat{f}_{a,b}(X_1)$, $\hat{f}_{a,b}(X_{-1})$, (b-1,b), (a,a+1), and $|\hat{f}'_{a,b}| \ge \min\{1/a^2, 1/b^2\} > 1$ on X. Zweimüller [23] showed that any one-dimensional map for which conditions (A), (F), (U) hold is exact and satisfies Rychlik's conditions described in [18], hence its natural extension map is Bernoulli.

We now analyze the case $b \ge 1$. Let K > 0 be the smallest integer such that $b(a + 1)^K < 1$. We will show that there exists $\gamma > 1$ such that, for every $x \in \bigcap_{i=0}^K \hat{f}_{a,b}^{-i}(X)$, some iterate $\hat{f}_{a,b}^n(x)$ with $n \le K + 1$ is expanding, i.e. $|(\hat{f}_{a,b}^n)'(x)| \ge \gamma$. (For the rest of the proof, we simplify the notation and let \hat{f} denote the map $\hat{f}_{a,b}$.) Notice that if $x \in \bigcap_{i=0}^{n-1} \hat{f}^{-i}(X)$, then \hat{f}^n is differentiable at x and

$$\frac{d}{dx}\hat{f}^n(x) = \frac{1}{(x\,\hat{f}(x)\cdots\hat{f}^{n-1}(x))^2}.$$

Assume that ab > -1. We look at the following cases:

- (i) If a < x < 0, then $b 1 \le \hat{f}(x) \le b$, and $|x \hat{f}(x)| \le |ab| < 1$.
- (ii) If 0 < x < b, then $a \le \hat{f}(x) \le a + 1$. Let K be such that $b(a+1)^K < 1$. Then either there exists $1 \le n \le K$ such that $0 < \hat{f}^i(x) < a + 1$ for $i = 1, 2, \ldots, n 1$ and $a < \hat{f}^n(x) < 0$, or $0 < \hat{f}^i(x) < a + 1$ for $i = 1, 2, \ldots, K$. In the former case we have that

$$|x \hat{f}(x) \cdots \hat{f}^{n}(x)| \le |ab(a+1)^{n-1}| < 1,$$
 (6.6)

while in the latter case

$$|x\hat{f}(x)\cdots\hat{f}^{K}(x)| \le |b(a+1)^{K}| < 1.$$
 (6.7)

In the case ab = -1, let τ , $\epsilon > 0$ be sufficiently small such that

$$b < -1/(a + \tau) < b + 1$$
 and $a - 1 < -1/(b - \epsilon) < a$.

We again consider two cases:

- (i) If $a < x < a + \tau$, then $b 1 < \hat{f}(x) < -1/(a + \tau)$, and $|x \hat{f}(x)| \le |a/(a + \tau)| < 1$. If $a + \tau \le x < 0$, then $|x \hat{f}(x)| \le |b(a + \tau)| < 1$.
- (ii) If $b \epsilon < x < b$, then $0 < \hat{f}(x) < a + 1$ and we have either (6.6) with $n \ge 2$ or (6.7). If $0 < x \le b \epsilon$, then we have (6.6) or (6.7) where b is replaced by $b \epsilon$.

In conclusion, there exists a constant $\gamma > 1$ such that for every $x \in \bigcap_{i=0}^K \hat{f}_{a,b}^{-i}(X)$ some iterate $\hat{f}_{a,b}^n(x)$ with $n \le K+1$ satisfies the condition $|(\hat{f}_{a,b}^n)'(x)| \ge \gamma$. This implies that the iterate $\hat{f}_{a,b}^N$, with N = (K+1)!, is uniformly expanding, i.e. it satisfies property (U). Since properties (A) and (F) are automatically satisfied by any iterate of $\hat{f}_{a,b}$ (see [23]), we have that $\hat{F}_{a,b}^N$ is Bernoulli. Using one of Ornstein's results [17, Theorem 4, p. 39], it follows that $\hat{F}_{a,b}$ is Bernoulli.

The next result gives a formula for the measure-theoretic entropy of $(\hat{F}_{a,b}, \nu_{a,b})$.

THEOREM 6.2. The measure-theoretic entropy of $(\hat{F}_{a,b}, \nu_{a,b})$ is given by

$$h_{\nu_{a,b}}(\hat{F}_{a,b}) = \frac{1}{K_{a,b}} \frac{\pi^2}{3}.$$
 (6.8)

Proof. To compute the entropy of this two-dimensional map, we use Abramov's formula [1],

$$h_{\tilde{m}}(\{\phi^t\}) = \frac{h_{\rho_{a,b}}(R_{a,b})}{\int_{\Lambda_{a,b}} g_{a,b} d\rho_{a,b}},$$

where \tilde{m} is the normalized Liouville measure $d\tilde{m} = dm/m(SM)$. It is well known that $m(SM) = \pi^2/3$ (see [3]) and $h_{\tilde{m}}(\{\phi^t\}) = 1$ (see [20]). The measure $d\tilde{m}$ can be represented by the Ambrose–Kakutani theorem [5] as a smooth probability measure on the space $\Lambda_{a,b}^{g_{a,b}}$,

$$d\tilde{m} = \frac{d\rho_{a,b} dt}{\int_{\Lambda_{a,b}} g_{a,b} d\rho_{a,b}},\tag{6.9}$$

where $d\rho_{a,b}$ is the probability measure on the cross-section $\Lambda_{a,b}$ given by (6.4). This implies that

$$d\tilde{m} = \frac{d\rho \ dt}{K_{a,b} \int_{\Lambda_{a,b}} g_{a,b} \ d\rho_{a,b}} = \frac{dm}{K_{a,b} \int_{\Lambda_{a,b}} g_{a,b} \ d\rho_{a,b}}.$$

Therefore $K_{a,b} \int_{\Lambda_{a,b}} g_{a,b} d\rho_{a,b} = m(SM) = \pi^2/3$ and

$$h_{\nu_{a,b}}(\hat{F}_{a,b}) = h_{\rho_{a,b}}(R_{a,b}) = \int_{\Lambda_{a,b}} g_{a,b} \, d\rho_{a,b} = \frac{1}{K_{a,b}} \frac{\pi^2}{3}.$$

Since $(\hat{F}_{a,b}, \nu_{a,b})$ is the natural extension of $(\hat{f}_{a,b}, \mu_{a,b})$, the measure-theoretic entropies of the two systems coincide, hence

$$h_{\mu_{a,b}}(\hat{f}_{a,b}) = \frac{1}{K_{a,b}} \frac{\pi^2}{3} . \tag{6.10}$$

As an immediate consequence of the above entropy formula we derive a growth rate relation for the denominators of the partial quotients p_n/q_n of (a, b)-continued fraction expansions, similar to the classical cases.

PROPOSITION 6.3. Let $\{q_n(x)\}$ be the sequence of the denominators of the partial quotients p_n/q_n associated to the (a,b)-continued fraction expansion of $x \in [a,b)$. Then

$$\lim_{n \to \infty} \frac{\log q_n(x)}{n} = \frac{1}{2} h_{\mu_{a,b}}(\hat{f}_{a,b}) = \frac{1}{K_{a,b}} \frac{\pi^2}{6} \quad \text{for a.e. } x.$$
 (6.11)

Proof. The proof is similar to the classical case: using Birkhoff's ergodic theorem, we have

$$\lim_{n \to \infty} \frac{\log q_n(x)}{n} = -\int_a^b \log |x| \, d\mu_{a,b}.$$

At the same time, Rokhlin's formula tells us that

$$h_{\mu_{a,b}}(\hat{f}_{a,b}) = \int_a^b \log|\hat{f}'_{a,b}| \, d\mu_{a,b} = -2 \int_a^b \log|x| \, d\mu_{a,b},$$

hence the conclusion.

7. Some explicit formulas for the invariant measure $\mu_{a,b}$

In order to obtain explicit formulas for $\mu_{a,b}$ and $h_{\mu_{a,b}}(\hat{f}_{a,b})$, one obviously needs an explicit description of the domain $D_{a,b}$. In [16] we describe an algorithmic approach

for finding the boundaries of $D_{a,b}$ for all parameter pairs (a, b) outside of a negligible exceptional parameter set \mathcal{E} . Let us point out that the set $D_{a,b}$ may have an arbitrary large number of horizontal boundary segments. The qualitative structure of $D_{a,b}$ is given by the cycle properties of a and b. This structure remains unchanged for all pairs (a, b) having cycles with similar combinatorial complexity. For a large part of the parameter set the cycle descriptions are relatively simple (see [16, §4]) and we discuss them here.

In what follows, we focus our attention on the situation $-1 \le a \le 0 \le b \le 1$, and due to the symmetry of the parameter set with respect to the parameter line a = -b we assume that $a \le -b$.

We treat the case $1 \le -1/a \le b+1$ and $a \le -1/b+m \le a+1$ (for some $m \ge 1$). The coordinates of the corners of the boundary segments in the upper region $D_{a,b} \cap \{(u, w) \mid u < 0, a \le w \le b\}$ are given by

$$(-2, b-1), \left(-\frac{3}{2}, T^{-2}S(b-1)\right), \dots, \left(-\frac{m+1}{m}, (T^{-2}S)^{(m-1)}(b-1)\right),$$

 $\left(-1, -\frac{1}{a} - 1\right),$

while the corners of the boundary segments in the lower region $D_{a,b} \cap \{(u, w) \mid u > 0, a \le w \le b\}$ are given by

$$\left(m, -\frac{1}{b} + m\right), (m+1, a+1).$$

Therefore the set $\hat{\Lambda}_{a,b}$ is given by

$$\hat{\Lambda}_{a,b} = \bigcup_{p=1}^{m-1} [(T^{-2}S)^{p-1}(b-1), (T^{-2}S)^{p}(b-1)] \times \left[0, \frac{p}{p+1}\right]$$

$$\cup \left[(T^{-2}S)^{m-1}(b-1), -\frac{1}{a}-1\right] \times \left[0, \frac{m}{m+1}\right] \cup \left[-\frac{1}{a}-1, b\right] \times [0, 1]$$

$$\cup \left[a, -\frac{1}{b}+m\right] \times \left[-\frac{1}{m}, 0\right] \cup \left[-\frac{1}{b}+m, a+1\right] \times \left[-\frac{1}{m+1}, 0\right]. \quad (7.15)$$

Figure 8 shows a typical set $\hat{\Lambda}_{a,b}$ for this case.

THEOREM 7.1. If $1 \le -1/a \le b + 1$ and $a \le -(1/b) + m \le a + 1$, then

$$\mu_{a,b} = \frac{1}{K_{a,b}} h_{a,b}(x) dx,$$

where $K_{a,b} = \log[(m-a)(1+b)^{2-m}]$ and $h_{a,b}(x) = h_{a,b}^+(x) + h_{a,b}^-(x)$ with

$$h_{a,b}^{+}(x) = \begin{cases} \frac{1}{x + ((p+1)/p)} & \text{if } (T^{-2}S)^{p-1}(b-1) \le x < (T^{-2}S)^{p}(b-1), \\ p = 1, \dots, m-1, \\ \frac{1}{x + ((m+1)/m)} & \text{if } (T^{-2}S)^{m-1}(b-1) \le x < -\frac{1}{a} - 1, \\ \frac{1}{x+1} & \text{if } -\frac{1}{a} - 1 \le x < b \end{cases}$$

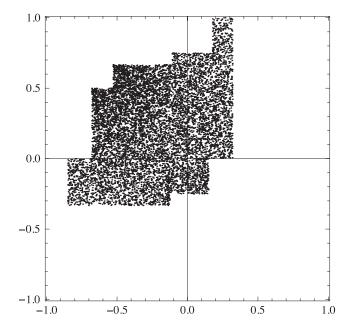


FIGURE 8. Typical domain $\hat{\Lambda}_{a,b}$ for the case studied.

and

$$h_{a,b}^{-}(x) = \begin{cases} \frac{1}{m-x} & \text{if } a \le x < -\frac{1}{b} + m, \\ \frac{1}{m+1-x} & \text{if } -\frac{1}{b} + m \le x < a+1. \end{cases}$$

Proof. The density formulas are obtained from the simple integration result

$$\int_{c}^{d} \frac{1}{(1+xy)^{2}} dy = -\frac{1}{x} \left(\frac{1}{1+dx} - \frac{1}{1+cx} \right) = \frac{d}{1+dx} - \frac{c}{1+cx}.$$
 (7.2)

For the density in the upper part of $\hat{\Lambda}_{a,b}$, $y \geq 0$, all integrals have the lower boundary c=0, hence the result of (7.2) becomes 1/(x+1/d). This gives the description of $h_{a,b}^+(x)$. For the density in the lower part of $\hat{\Lambda}_{a,b}$, $y \leq 0$, all integrals have the upper boundary d=0, hence the result -1/(-1/c-x) and the description of $h_{a,b}^-(x)$. By a somewhat tedious computation, we get

$$K_{a,b} = \int_{\Lambda_{a,b}} h_{a,b}(x) dx = \log[(m-a)(1+b)^{2-m}],$$

and this completes the proof.

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